

CHAPTER III

VELOCITY - DEPTH DISTRIBUTION

3.1 Computation of P Wave Velocity :

There are several methods of inverting a travel time curve to get a velocity distribution at different depths which have been adopted by various workers since the beginning of seismology.

Dowling and Nuttli (1964), studying for a low velocity channel in the upper mantle, used a modification of the following formulas (Bullen, 1963) -

$$\Delta = 2 \int_{r_m}^{r_0} \frac{p}{r} \frac{dr}{\sqrt{\{(r/v)^2 - p^2\}}} \quad 3.1a$$

and

$$T = 2 \int_{r_m}^{r_0} \frac{K}{v^2} \frac{dr}{\sqrt{\{(r/v)^2 - p^2\}}} \quad 3.1b$$

Considering the velocity function $v(r)$ in some constant steps, i.e. considering the earth to be consisted of some concentric spherical shells within which the ray path remains straight, the expression for the epicentral distance and travel time as given in the set of equations 3.1 above were approximated by them as a summation of two series

as given below.

$$\Delta = 2 \sum \frac{p (r_i - r_{i+1})}{r_i \sqrt{(r_i/v_i)^2 - p^2}} \quad 3.2a$$

and

$$T = 2 \sum \frac{r_i (r_i - r_{i+1})}{v_i^2 \sqrt{(r_i/v_i)^2 - p^2}} \quad 3.2b$$

fixing the ray parameter p by determining the specific angle of incidence, i_f , at the focus by the relation,

$$p = (r_f/v_f) \sin i_f \quad 3.3$$

Taking comparatively a larger number of layers, and assigning values of velocity from a simple velocity model and adjusting them particularly for suspected anomalous velocity regions, they could identify the possible existence of a low velocity channel from the shadow zones as evidenced from their calculated $T - \Delta$ curves.

Carder (1964) in his study of nuclear explosions in the Central Pacific determined P wave velocities using the basic principles of ray theory, within the epicentral distance range of $3^\circ - 102^\circ$ and assuming the mantle to be divided in eight 'constant velocity' concentric layers on the basis of straight line segments fitted to his $T - \Delta$ curve. Kaila et al. (1968b), Biswas and Bhattacharyya (1974) also, among others, adopted similar techniques with some modifications for studying P wave velocity structure in the upper

mantle under Indian sub-continent and near Alaska respectively.

Basanova et al. (1974; 1976) used the $t_0 - p$ method for inversion of the travel time taking into account the delay time function $\tau(p)$ along with the consideration of the apparent slowness characterizing a velocity structure of a strictly vertically heterogeneous medium, later on used again in a geometrically modified form by Germany and Orcutt (1979). Monte-Carlo method of inversion of (T, Δ) data is also used by many (Press, 1968; 1970; Wiggins, 1969) for similar purposes.

The methods have their own advantages, specially when applied to a mixture of data, but for the continuous inversion of the travel time data alone, they appear to be unnecessarily complicated. Comparison between various methods have also been made (Wiggins et al., 1973), although it is not obvious how good the comparisons should be, "as methods other than Herglotz - Wiecherdt inversion gives averaged uncertainties only" (Davis and Chapman, 1975).

3.2 Herglotz - Wiecherdt - Bateman Integral Method :

This integral method is chiefly a refraction method following a principle of optical refraction technique and transformed from apparent surface velocity from the travel time data to a velocity depth inversion by graphical integration method (Byerly, 1956).

Hamilton's equation for wave propagation is,

$$\left(\frac{\partial T}{\partial x}\right)^2 + \left(\frac{\partial T}{\partial y}\right)^2 + \left(\frac{\partial T}{\partial z}\right)^2 = \frac{1}{v^2} \quad 3.4$$

If the speed v is a function of z alone the rays with each lie in a vertical plane containing the source and the observation point, one can take X axis in this direction $\partial T/\partial x$ becomes the gradient of scalar time field created by the wave arrived at the Earth's surface. (Throughout this treatment, v will denote the true velocity and V will denote apparent velocity.)

Equation 3.4 can now be written as,

$$\left(\frac{\partial T}{\partial x}\right)^2 = \frac{1}{v^2} - \left(\frac{\partial T}{\partial z}\right)^2 \quad 3.5$$

A solution to the above equation may be obtained by setting p ($= \partial T/\partial x$) as the ray parameter.

Equation 3.5 is to be used for the areas of the Earth's surface small enough to neglect its curvature. In this case from Snell's - Descartes' law,

$$p = \frac{\sin i}{v} = \frac{\sin i_0}{v_0} = \frac{1}{V} \quad 3.6$$

where V ($= 1/p$) is the apparent speed at the surface, where the rays emerge and is obtained from the slope of the travel time curve. The angle i is between the ray and the

vertical and the suffix (o) denotes the parameters at the surface of the Earth.

Now applying equation 3.5 for the larger areas on the surface of the Earth and changing to polar co-ordinates with the centre of the Earth as the origin, the equation can be written as,

$$\frac{r^2}{v^2} - r^2 \left(\frac{\partial T}{\partial r} \right)^2 = \left(\frac{\partial T}{\partial \theta} \right)^2 \quad 3.7$$

Here it is assumed that v is a function of r alone and the solution is obtained by setting $p (= \partial T / \partial \theta)$. It may then be shown that,

$$p = \frac{r \sin i}{v} = \frac{r_0 \sin i_0}{v_0} = \frac{r_0}{v_0} \quad 3.8$$

where i is the angle made by the ray with the vertical, i.e. the radius of the Earth.

It follows then,

$$dT = \frac{\partial T}{\partial \theta} d\theta + \frac{\partial T}{\partial r} dr \quad 3.9$$

From 3.7, using p ,

$$T = \int p d\theta \pm \int \frac{dr}{r} \sqrt{(r/v)^2 - p^2} \quad 3.10$$

To get the equation of the ray,

$$\frac{\partial T}{\partial p} = \int d\theta + \int \frac{r dr}{r \sqrt{(r/v)^2 - p^2}} \quad 3.11$$

one may now integrate from the bottom of the ray (ref. to the Figures 10a and 10b) to the surface of the Earth and obtain,

$$\Delta = 2 \int_r^{r_0} \frac{r dr}{\sqrt{(r/v)^2 - p^2}} \quad 3.12$$

The difficulty in equation 3.12 is that, since it is desired to get v as a function of r , the integration cannot be performed along the ray, unless it is known which function will give v as a function of r . The travel time gives p as a function of Δ ($\equiv \theta$).

Herglotz in 1907, using Abel's transformation (Macelwane, 1932; Byerly, 1942) transformed the relation 3.12 to,

$$\ln (r_0/r_m) = \frac{1}{\pi} \int_0^{\Delta_m} \cosh^{-1} (v_m/v) d\Delta \quad 3.13a$$

$$= \frac{1}{\pi} \int_0^{\Delta_m} \cosh^{-1} (p/p_m) d\Delta \quad 3.13b$$

Now the integration is no longer along a ray as in equation 3.12 but along the Earth's surface; that is along the travel time curve (ref. to Figures 11a and 11b).

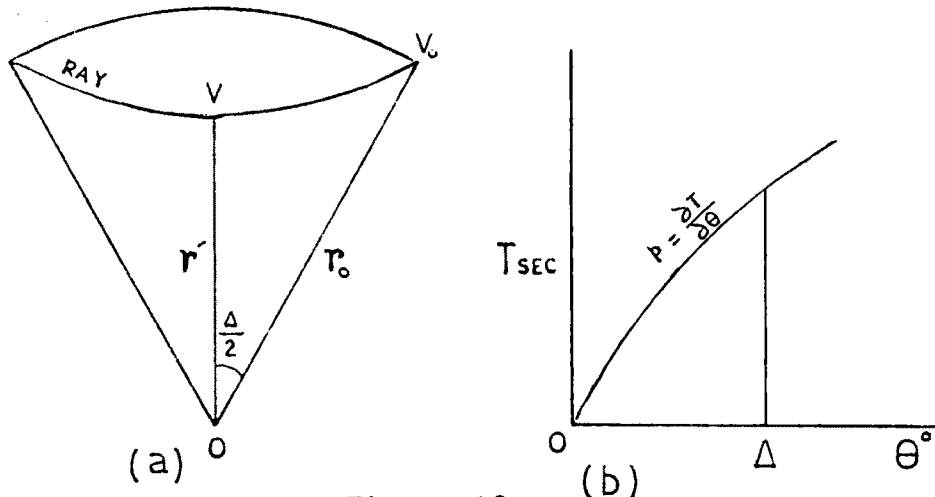


Figure 10

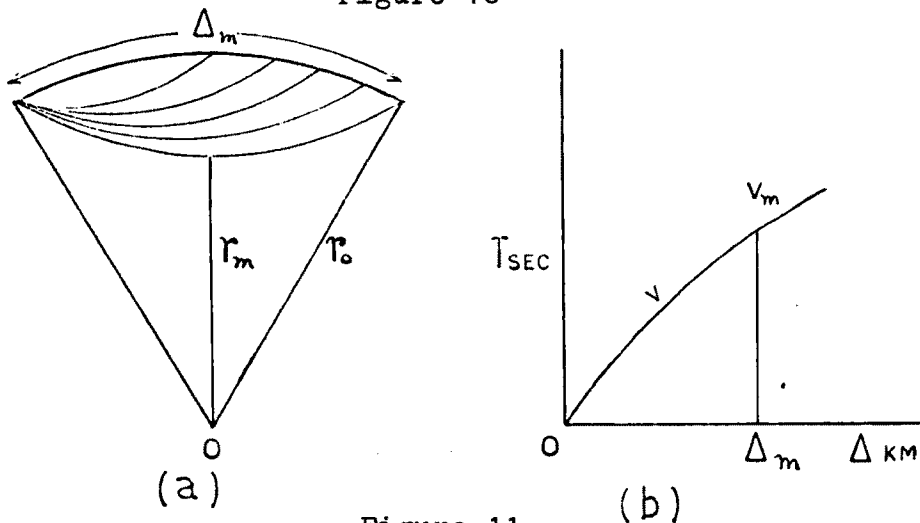


Figure 11

FIGURE 10 : (a) Ray diagram & (b) corresponding travel time curve

FIGURE 11 : (a) Ray diagram & (b) travel time curve illustrating Herglotz transformation

The apparent surface velocity $V (\equiv 1/p)$ is obtained from the slope of the travel time curve as a function of arc distance. Thus the equation 3.13b can be integrated graphically or numerically by selecting a value of p_m (or V_m) corresponding to a certain Δ_m and the integration is done from the surface of the Earth to evaluate the corresponding r_m , the vertex at which the ray with the ray parameter p_m bottoms. At this corresponding depth the velocity of the wave is,

$$v_m = \frac{r_m}{r_0} \cdot V_m \quad 3.14$$

The trouble in Herglotz method is that at the end of each integration, velocity at one depth only is obtained. To calculate the velocity again at a different depth the Earth is to be stripped and the epicentral distance Δ is to be properly reduced to bring the rays to the next selected depth again to a single point on the new surface of the stripped Earth.

This method is again applicable only when the following conditions are satisfied during the process of integration, otherwise the method fails as the argument of the integration becomes imaginary. The conditions are (Macelwane, 1932),

(i) dv/dr is -ve (velocity increases with depth) and

(ii) dv/dr is +ve but numerically less than v/r .

If at a certain level this criterion is not met and if the extent of the zone is comparatively not very wide, it is customary to assume that the region is a discontinuity for all practical purposes and the integration is performed upto that level and then the Earth is stripped to start the integration again to a new depth after proper reduction of the parameters involved in the integration.

3.3 Stripping of the Earth and Reduction of Parameters :

In the present study since the $p - \Delta$ profiles are fitted with polynomial segments in-between the assumed discontinuity levels the relation 3.13 and 3.14 are used in the following form for a selected level j .

$$\ln r_j = \ln r - \frac{1}{\pi} \int_0^{\Delta_j} \cosh^{-1} (p_{\Delta} / p_j) d\Delta_j \quad 3.15$$

and
$$v_j = r_j / (r_0 p_j) \quad 3.16$$

Here,
 r_j = radius of the Earth at the vertex of the ray bottoming at a selected depth,
 r = radius of the stripped Earth
 p_j = ray parameter of the ray emerging at the selected epicentral distance Δ_j

p_{Δ} = ray parameter of the ray emerging at an epicentral distance Δ within $0 - \Delta_j$

and, r_0 = radius of the Earth.

Now referring to Figure 12, let

\bar{v}_A = true velocity at the surface of the Earth corresponding to arc distance Δ measured at the surface of the Earth ($R = r_0$),

\bar{v}' = true velocity at the surface of the strip-ped Earth,

\bar{v}'_r = true velocity at the vertex of the ray ($R = r_m$)

V' = apparent velocity from T - Δ curve of a ray

Now from the figure - using Snell's law,

$$\frac{\sin \theta}{\sin \phi} = \frac{\bar{v}'_r}{\bar{v}'} \quad 3.17$$

and also,

$$\sin \theta = \frac{\bar{v}'_r}{V'} \quad (\text{Macelwane, 1932}) \quad 3.18$$

Hence from 3.17 and 3.18

$$\begin{aligned} \sin \phi &= (\bar{v}' / \bar{v}'_r) \sin \theta \\ &= \bar{v}' / V' \\ &= \bar{v}' \cdot p_{\Delta} \end{aligned} \quad 3.19$$

where, $p_{\Delta} = 1/V' = \frac{180}{\pi r_0} \frac{dt}{d\Delta} \text{ sec/Km}$

Hence, $\phi = \sin^{-1} (C_1 \cdot p_{\Delta}) \quad 3.20$

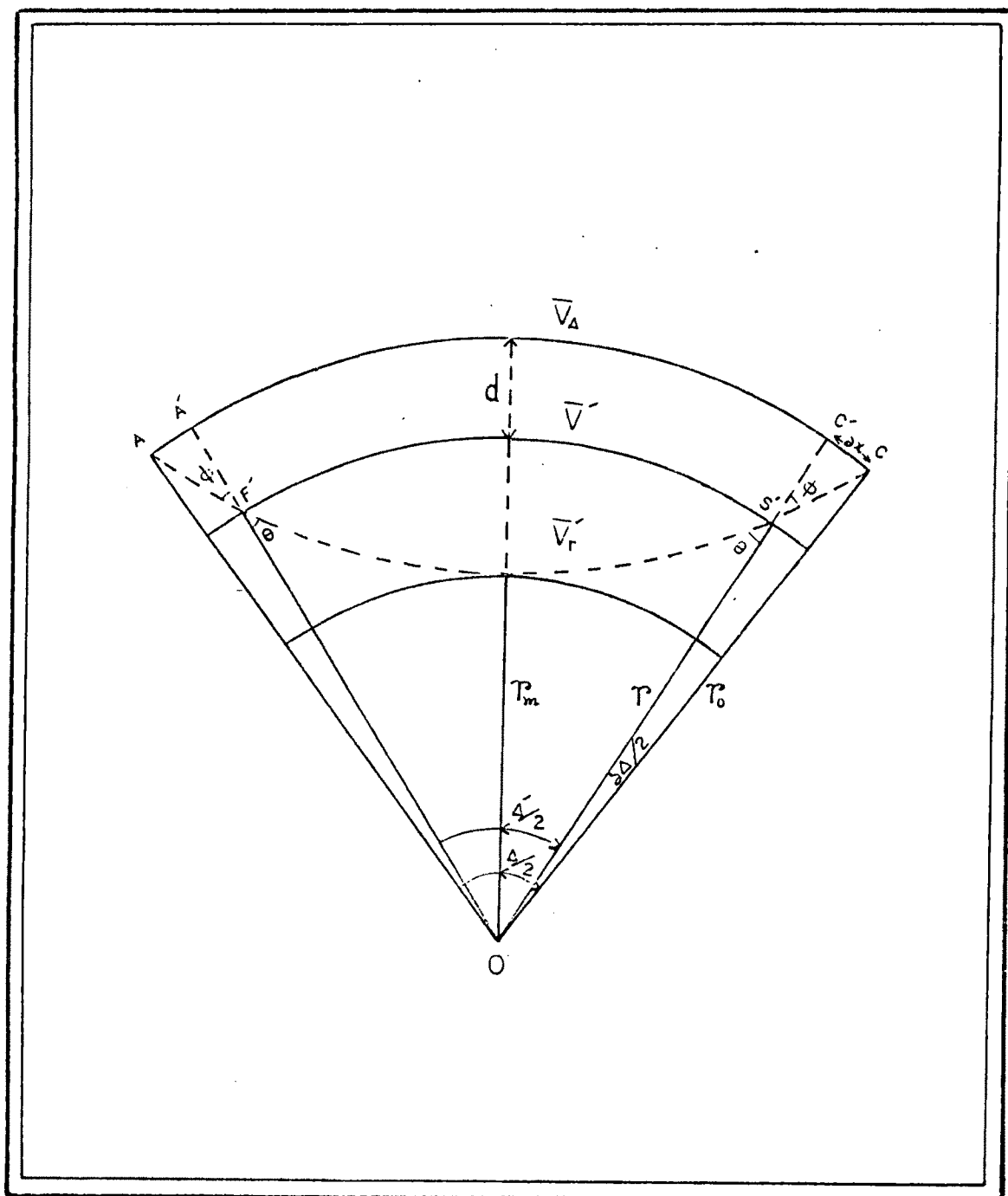


FIGURE 12 : Stripping of the Earth and reduction of parameters.

where, $C_1 = (180/\pi r_0) \cdot \bar{v}'$, is a constant for a particular layer selected after fixing the value of \bar{v}' for that layer (Appendix III).

If the upper layer is stripped the ray is assumed to originate at F' and emerge at S' on the surface of the stripped Earth covering a reduced epicentral distance Δ' corresponding to the surface of the Earth. From the Figure 12,

$$\Delta' = \Delta - 2\partial\Delta/2 \quad 3.21$$

It is clear from the figure that on the surface of the Earth this distance is

$$A'C' = AC - 2\partial x$$

$$\text{Now, } \partial x = d \tan \phi = r_0 \partial\Delta/2 \quad 3.22$$

$$\begin{aligned} \text{Hence } \partial\Delta &= 2d \tan \phi / r_0 \\ &= (360/\pi r_0) d \tan \phi \\ &= C_2 \cdot \tan \phi \end{aligned} \quad 3.23$$

Here, $C_2 = (360/\pi r_0) \cdot d =$ a constant for the assumed layer at depth d from the surface .

If Δ'_1 ($i = 1, 2, 3, \dots$) be the reduced arc distances on the surface of the stripped Earth corresponding to the arc distances Δ_1 on the surface of the true Earth then the corresponding decrements are,

$$\partial\Delta_1 = \Delta_1 - \Delta'_1 \quad 3.24$$

Then the increments of Δ along the surface of the Earth will be

$$d \Delta_j = \delta \Delta_i - \delta \Delta_{i-1} \quad 3.25$$

For computing the velocity distribution with depth, the integration in equation 3.15 has been made by using trapezoidal rule and by using the relations 3.16 to 3.25 in a computer programme

Since in this study the velocity distribution is restricted to the lower part of the mantle below the maximum depth of penetration of a ray emerging at about 30° epicentral distance a standard (Herrin, 1968) P wave velocity model has been used for calculation of velocity v (Appendix IIIA) for the upper part of the mantle stripped at 750 Km depth initially to start the integration and reduction of Δ and evaluation of the constants.

Here quite a few strippings had to be undertaken to compute the velocity - depth distribution bounded by the anomalous zones in each profile and in all the cases the upper part of the Earth had been stripped and the average velocity was calculated (Appendix III); the starting level being taken from the end point of the $p - \Delta$ profile (segment) of the previous upper layer.

The velocities within each segment have been calculated at an interval of 0.5° arc distance and

finally the average velocity is calculated using the raw Herglotz - Wiechert values for each profile, by polynomial method of interpolation at 50 Km depth interval. The computed Herglotz - Wiechert values and the interpolated average velocity - depth distributions for all the profiles are shown in Figures 13a and 13b. Herrin's (1968) standard surface focus average velocity distribution is also shown in each profile for comparison. The values of the interpolated average velocity - depth distribution at 50Km depth interval are presented in Table V.

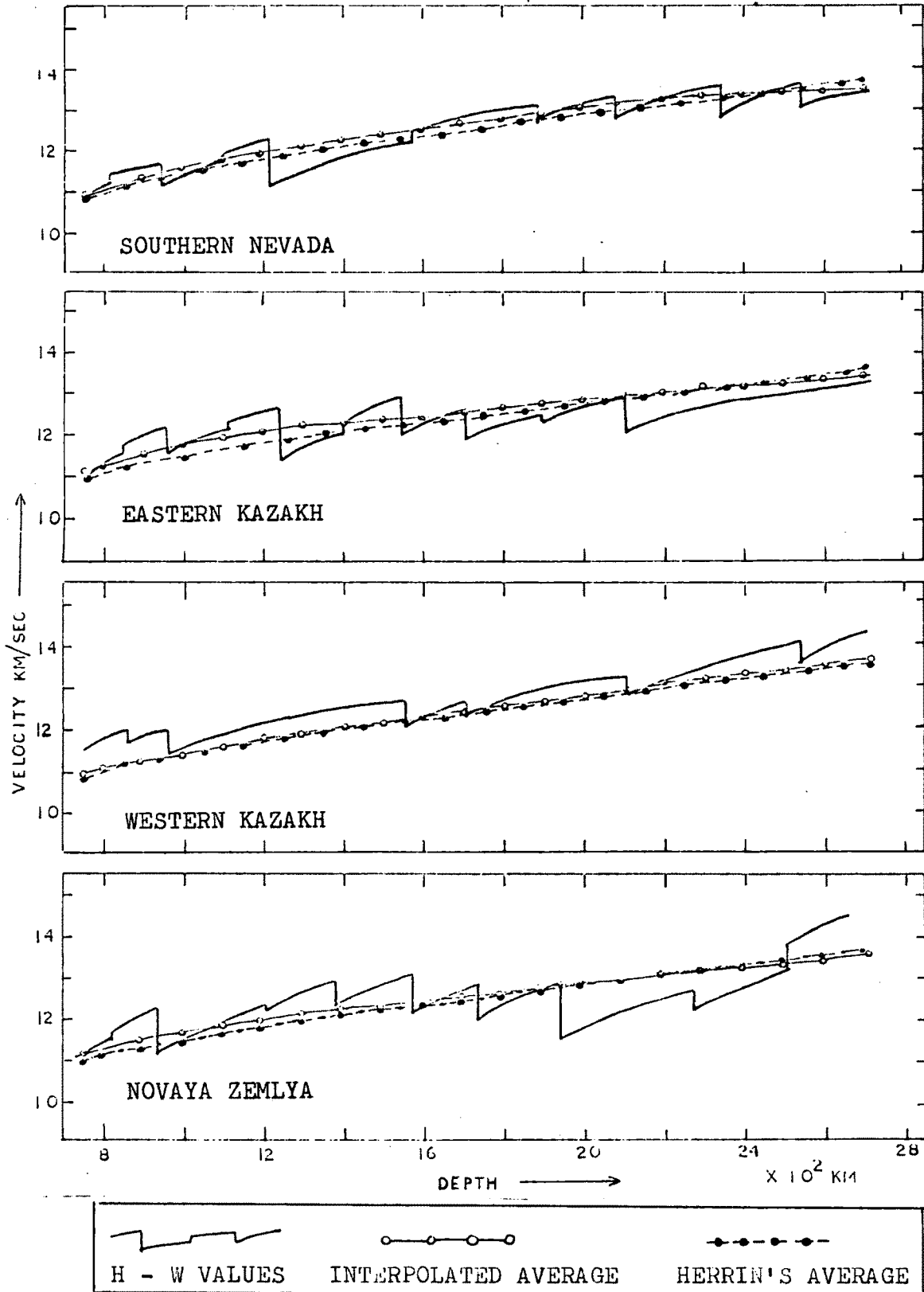


FIGURE 13a : Velocity - depth distribution for the nuclear explosion profiles, compared with Herrin's(1968) values for surface focus event.

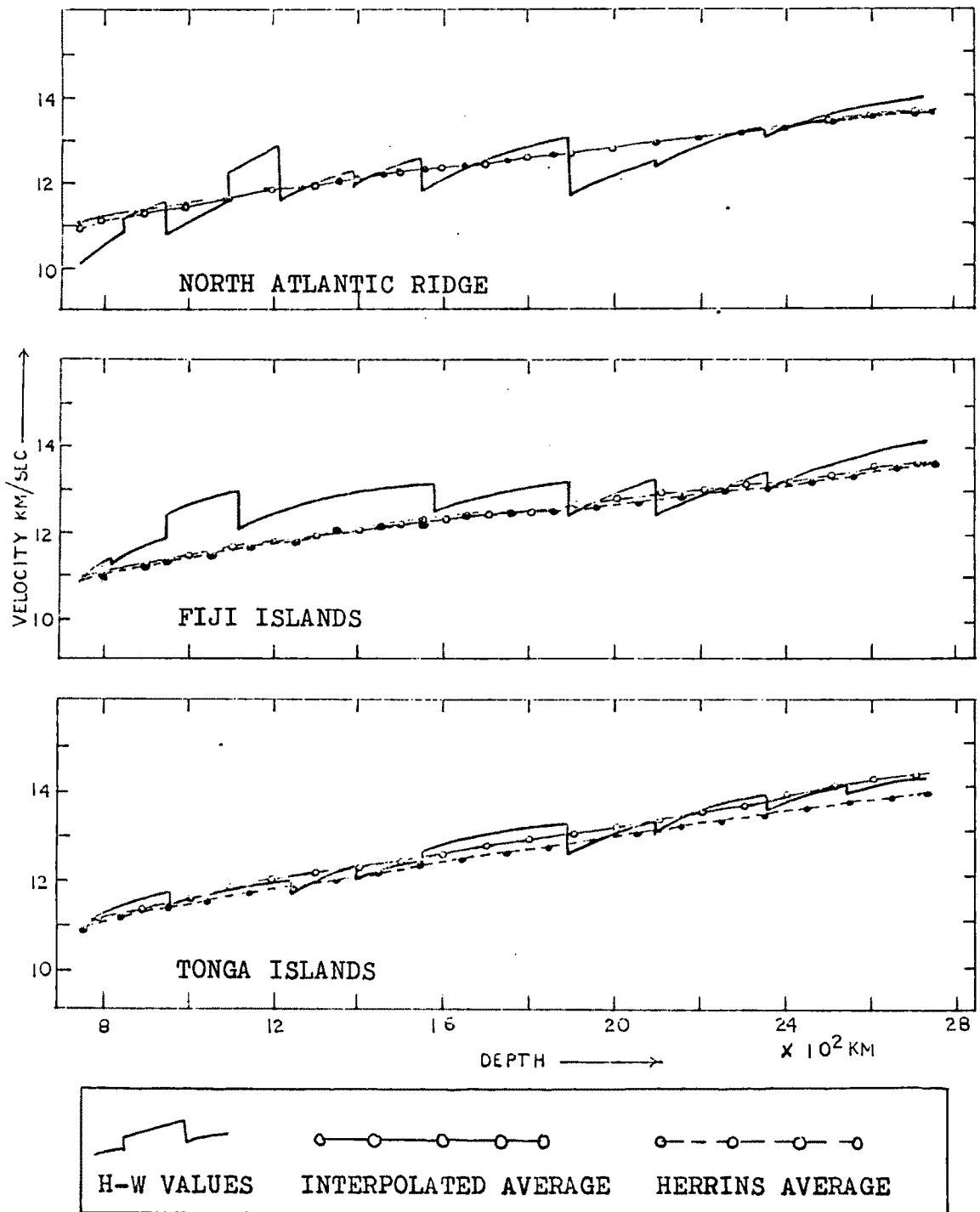


FIGURE 13b : Velocity¹ - depth distribution for the earthquake profiles, compared with Herrin's (1968) values for surface focus event.

TABLE V : Average velocity - depth distribution from Herglotz - Wiechert values for
each profile

| Depth Km | P wave velocity in Km/sec. for - | | | | | | | |
|-------------|----------------------------------|-------------------|-------------------|------------------|----------------------------|-----------------|------------------|--|
| | Southern Nevada | Eastern Kazakh | Western Kazakh | Novaya Zemlya | North Atlantic Ridge | Fiji Islands | Tonga Islands | |
| 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | |
| 750 | 11.0545 | 11.3271 | 11.0570 | 11.0860 | 10.9752 | 10.9827 | 11.0923 | |
| 800 | 11.1695 | 11.4612 | 11.1498 | 11.3847 | 11.0586 | 11.0666 | 11.1804 | |
| 850 | 11.2810 | 11.5823 | 11.2402 | 11.6471 | 11.1411 | 11.1496 | 11.2674 | |
| 900 | 11.3895 | 11.6512 | 11.3282 | 11.7642 | 11.2227 | 11.2617 | 11.3545 | |
| 950 | 11.4949 | 11.7886 | 11.4140 | 12.0642 | 11.3034 | 11.3829 | 11.4386 | |
| 1000 | 11.5973 | 11.8751 | 11.4976 | 12.2243 | 11.3832 | 11.4332 | 11.5328 | |
| 1050 | 11.6968 | 11.9515 | 11.5791 | 12.3551 | 11.4621 | 11.4926 | 11.6063 | |
| 1100 | 11.7935 | 12.0183 | 11.6585 | 12.4588 | 11.5400 | 11.5810 | 11.6890 | |
| 1150 | 11.8863 | 12.0764 | 11.7361 | 12.5377 | 11.6171 | 11.6586 | 11.7710 | |
| 1200 | 11.9776 | 12.1264 | 11.8117 | 12.5941 | 11.6932 | 11.7152 | 11.8525 | |
| 1250 | 12.0654 | 12.1690 | 11.8851 | 12.6301 | 11.7684 | 11.8009 | 11.9335 | |

continued ...

TABLE V : continued.

| 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 |
|------|---------|---------|---------|---------|---------|---------|---------|
| 1300 | 12.1505 | 12.2049 | 11.9578 | 12.6481 | 11.8428 | 11.8558 | 12.0141 |
| 1350 | 12.2327 | 12.2347 | 12.0284 | 12.6502 | 11.9162 | 11.9497 | 12.0942 |
| 1400 | 12.3122 | 12.2592 | 12.0974 | 12.6387 | 11.9887 | 12.0827 | 12.1741 |
| 1450 | 12.3890 | 12.2790 | 12.1650 | 12.6158 | 12.0603 | 12.1648 | 12.2537 |
| 1500 | 12.4630 | 12.2949 | 12.2313 | 12.5838 | 12.1310 | 12.1960 | 12.3332 |
| 1550 | 12.5344 | 12.3075 | 12.2962 | 12.5449 | 12.2007 | 12.2162 | 12.4126 |
| 1600 | 12.6031 | 12.3175 | 12.3600 | 12.5014 | 12.2696 | 12.2856 | 12.4920 |
| 1650 | 12.6692 | 12.3256 | 12.4226 | 12.4555 | 12.3375 | 12.3540 | 12.5714 |
| 1700 | 12.7328 | 12.3325 | 12.4841 | 12.4094 | 12.4046 | 12.4216 | 12.6509 |
| 1750 | 12.7938 | 12.3388 | 12.5447 | 12.3654 | 12.4707 | 12.4882 | 12.7307 |
| 1800 | 12.8522 | 12.3454 | 12.6045 | 12.3257 | 12.5360 | 12.5540 | 12.8107 |
| 1850 | 12.9082 | 12.3527 | 12.6634 | 12.2926 | 12.6003 | 12.6188 | 12.8910 |
| 1900 | 12.9617 | 12.3617 | 12.7216 | 12.2683 | 12.6637 | 12.6827 | 12.9717 |
| 1950 | 13.0128 | 12.3728 | 12.7792 | 12.2551 | 12.7262 | 12.7457 | 13.0529 |
| 2000 | 13.0614 | 12.3869 | 12.8362 | 12.2551 | 12.7878 | 12.8078 | 13.1347 |
| 2050 | 13.1077 | 12.4045 | 12.8927 | 12.2707 | 12.8485 | 12.8690 | 13.2171 |
| 2100 | 13.1516 | 12.4149 | 12.9489 | 12.3041 | 12.9083 | 12.9293 | 13.3001 |

continued ...

TABLE V : continued

| 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 |
|------|---------|---------|---------|---------|---------|---------|---------|
| 2150 | 13.1933 | 12.4393 | 13.0047 | 12.3575 | 12.9671 | 12.9886 | 13.3870 |
| 2200 | 13.2326 | 12.4861 | 13.0604 | 12.4331 | 13.0251 | 13.0471 | 13.4686 |
| 2250 | 13.2697 | 12.5250 | 13.1158 | 12.5332 | 13.0821 | 13.1046 | 13.5113 |
| 2300 | 13.3045 | 12.5711 | 13.1712 | 12.6601 | 13.1383 | 13.1613 | 13.6407 |
| 2350 | 13.3372 | 12.6248 | 13.2267 | 12.8160 | 13.1935 | 13.2170 | 13.7282 |
| 2400 | 13.3677 | 12.6870 | 13.2822 | 13.0031 | 13.2478 | 13.2718 | 13.8169 |
| 2450 | 13.3960 | 12.7583 | 13.3379 | 13.2236 | 13.3012 | 13.3257 | 13.9067 |
| 2500 | 13.4223 | 12.8394 | 13.3938 | 13.4799 | 13.3538 | 13.3788 | 13.9978 |
| 2550 | 13.4464 | 12.9310 | 13.4501 | 13.4840 | 13.4050 | 13.4308 | 14.0902 |
| 2600 | 13.4685 | 13.0338 | 13.5068 | 13.5484 | 13.4560 | 13.4820 | 14.1170 |
| 2650 | 13.4886 | 13.1484 | 13.5641 | 13.5952 | 13.5058 | 13.5323 | 14.1380 |
| 2700 | 13.5067 | 13.2756 | 13.6219 | 13.6568 | 13.5547 | 13.5817 | 14.1640 |
| 2750 | 13.5530 | 13.4160 | 13.6804 | 13.6971 | 13.6026 | 13.6301 | 14.1810 |