CHAPTER - 4

ON TWO - FLUID BLOOD FLOW
THROUGH STENOSED ARTERY WITH
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*The universe is strange,*
*it is stranger than you think,*
*it may be stranger than you can think.*
- Aldous Huxley

4.1 Introduction

The frequently occurring cardiovascular disease, arteriosclerosis or stenosis, responsible for many of the diseases is the unnatural and abnormal growth that develops at various locations of the cardiovascular system under diseased conditions. Although, the etiology of the initiation of the disease (stenosis) is not well understood but it is well established that once the constriction has developed, it brings about the significant changes in the flow field (pressure distribution, wall shear stress, impedance, etc.). The flowing blood has been represented by a Newtonian, non-Newtonian, single or double-layered fluid by the investigators in the literature while discussing the flow through stenoses. It is well known that blood may be represented by a single-layered model in large vessel; however, the flow through the small arteries is known to be a two-layered. Bugliarello and Sevilla (1970) and Cokelet (1972) have shown experimentally that for blood flowing through small vessels, there is cell-free plasma (Newtonian viscous fluid) layer and a core region of suspension of all the erythrocytes. Srivastava (2007) concluded that the significance of the peripheral layer increases with decreasing blood vessel diameter. In addition, the endothelial walls are known to be highly permeable with ultra a microscopic pore through which filtration
occurs. Cholesterol is believed to increase the permeability of the arterial wall. Such increase in permeability results from dilated, damaged or inflamed arterial walls. In view of the discussion given above, the research reported in this chapter is therefore devoted to discuss the two-layered blood flow through an axisymmetric stenosis in an artery with permeable wall. The mathematical model considers the flowing blood as a two-layered Newtonian fluid, consisting of a core region (central layer) of suspension of all the erythrocytes assumed to be a Newtonian fluid, the viscosity of which may vary depending on the flow conditions and a peripheral region (outer layer) of another Newtonian fluid (plasma) of constant viscosity, in an artery with permeable wall.

4.2 Formulation of the problem

Consider the two-layered axisymmetric flow of blood through an axisymmetric stenosis, specified at the location as shown in Fig. (4.2.1) The geometry of the stenosis which is assumed to be manifested in the arterial wall segment is described as

![Diagram](image)

Figure: 4.2.1. Two-layered blood flow through an axisymmetric stenosis with permeable wall.
\begin{align}
\frac{R(z), R_1(z)}{R_0} &= (1, \beta) - \left(\frac{\delta, \delta_1}{2R_0}\right) \left\{1 + \cos\frac{2\pi}{L_0}\left(z - d - \frac{L_0}{2}\right)\right\}; \quad d \leq z \leq d + L_0, \quad (4.2.1) \\
&= (1, \beta); \quad \text{otherwise,} \quad (4.2.2)
\end{align}

where \( z \) is the axial coordinate, \((R, R_1) \equiv (R(z), R_1(z))\) are the radius of the (tube, interface) with constriction; \( R_0 \) is the radius of the normal (without stenosis) artery; \( L_0 \) is the stenosis length, \( L \) is the tube length and \( d \) indicates the location of stenosis, \( \beta \) is the ratio of the central core radius to the tube radius in the unobstructed region and \((\delta, \delta_1)\) are the maximum height of the stenosis and the bulging of the interface.

The flowing blood is assumed to be represented by a two-layered Newtonian fluid. The equations describing the laminar, steady, one-dimensional flow in the case of a mild stenosis \((\delta \ll R_0)\) are expressed \(\text{(Sharan and Popel, 2001)}\) as

\begin{align}
\frac{dp}{dz} &= \frac{\mu_p}{r} \frac{\partial}{\partial r} (r \frac{\partial}{\partial r}) u_p, \quad R_1(z) \leq r \leq R(z), \quad (4.2.3) \\
\frac{dp}{dz} &= \frac{\mu_c}{r} \frac{\partial}{\partial r} (r \frac{\partial}{\partial r}) u_c, \quad 0 \leq r \leq R_1(z), \quad (4.2.4)
\end{align}

where \( R_1(z) \) is the radius of the central layer, \((u_p, \mu_p)\) and \((u_c, \mu_c)\) are (velocity, viscosity) of fluid in the peripheral layer \((R_1(z) \leq r \leq R(z))\) and central layer \((0 \leq r \leq R_1(z))\), respectively; \( p \) is the pressure and \((r, z)\) are (radial, axial) coordinates in the two-dimensional cylindrical polar coordinate system.

The appropriate boundary conditions \(\text{(Beavers and Joseph, 1967)}\) for the present problem may be stated \(\text{(Srivastava et al., 2012)}\) as

\[ \frac{\partial u_c}{\partial r} = 0 \quad \text{at} \quad r = 0, \quad (4.2.5) \]
\[ u_p = u_c \text{ and } \mu_p \frac{\partial u_p}{\partial r} = \mu_c \frac{\partial u_c}{\partial r} \text{ at } r = R_1(z), \quad (4.2.6) \]

\[ u_p = u_B \text{ and } \frac{\partial u_p}{\partial r} = \frac{\alpha}{\sqrt{k}} (u_B - u_{\text{porous}}) \text{ at } r = R(z), \quad (4.2.7) \]

where \( u_{\text{porous}} = -\frac{k}{\mu_p} \frac{dp}{dz} \), is the velocity in the permeable boundary, \( u_B \) is the slip velocity, \( \mu_p \) is the plasma viscosity (fluid viscosity in the peripheral layer), \( k \) is Darcy number and \( \alpha \) (called the slip parameter) is a dimensionless quantity depending on the material parameters which characterize the structure of the permeable material within the boundary region.

### 4.3 Analysis

The solution of Eqs. (4.2.3) and (4.2.4) under the boundary conditions (4.2.5), (4.2.6) and (4.2.7), yields the expressions for velocity, \( u_p \) and \( u_c \) as

\[ u_p = \frac{-R_0^2}{4\mu_p} \frac{dp}{dz} \left\{ \left( \frac{R}{R_0} \right)^2 - \frac{r}{R_0} \left( \frac{R}{R_0} \right)^2 - 2 \left( \frac{R}{R_0} \right)^2 \frac{\sqrt{k}}{\alpha R_0} + 4 \frac{k}{R_0^2} \right\}, \quad (4.3.1) \]

\[ u_c = \frac{-R_0^2}{4\mu_p} \frac{dp}{dz} \left\{ \left( \frac{R}{R_0} \right)^2 - \mu \left( \frac{R}{R_0} \right)^2 - (1 - \mu) \left( \frac{R_1}{R_0} \right)^2 - 2 \left( \frac{R}{R_0} \right)^2 \frac{\sqrt{k}}{\alpha R_0} + 4 \frac{k}{R_0^2} \right\}, \quad (4.3.2) \]

with \( \mu = \frac{\mu_p}{\mu_c} \).

The volumetric flow rate, \( Q \) is now calculated as

\[ Q = 2\pi \left\{ \int_0^{R_1} \left[ \frac{r}{u_p} \frac{dr}{u_c} + \frac{r}{u_i} \frac{dr}{u_i} \right] \right\}, \]

\[ Q = -\frac{\pi R_0^4}{8\mu_p} \frac{dp}{dz} \left\{ \left( \frac{R}{R_0} \right)^4 - (1 - \mu) \left( \frac{R_1}{R_0} \right)^2 + \frac{8k}{R_0^2} \left( \frac{R}{R_0} \right)^2 - 4\sqrt{k} \left( \frac{R}{R_0} \right)^3 \right\}, \quad (4.3.3) \]

Following the argument \((Shukla et al., 1980; Srivastava et al., 2010c)\) that the total
of the fluxes across the two regions (peripheral and central), one derives the relations:

\[ R_1 = \beta R \quad \text{and} \quad \delta_1 = \alpha \delta \left(0 \leq \beta \leq 1\right). \]

An application of these relations into the Eq. (4.3.3), yields

\[
\frac{dp}{dz} = -\frac{8 \mu_p Q}{\pi R_0^4} \varphi(z),
\]

where

\[ \varphi(z) = \sqrt{1 - (1 - \mu)\beta^4 \left(R/R_0\right)^4 + 8k \left(R/R_0\right)^2/R_0^3 - 4\sqrt{k} \left(R/R_0\right)^2/\alpha R_0}. \]

The pressure drop, \( \Delta p = p \text{ at } z = 0, -p \text{ at } z = L \) across the stenosis in the tube of length, \( L \) is obtained as

\[
\Delta p = \int_0^L \left( -\frac{dp}{dz} \right) dz
= \frac{8 \mu_p Q}{\pi R_0^4} \left\{ \int_0^d \varphi(z) R/R_n=1 \, dz + \int_d^{d+L_n} \varphi(z) \, dz + \int_{d+L_n}^L \varphi(z) R/R_n=1 \, dz \right\},
\]

(4.3.5)

The analytical evaluation of the second integral on the right hand side of Eq. (4.3.5) is a formidable task and therefore shall be evaluated numerically, whereas the evaluation of first and third integrals are straightforward. Using now the definitions from the published literature \( \text{(Young, 1968; Srivastava et al., 2010c)} \), one derives the expressions for the impedance (flow resistance), \( \lambda \), the wall shear stress distribution in stenotic region, \( \tau_w \) and the shear stress at the stenosis throat, \( \tau_s \) in their non-dimensional form as

\[
\lambda = \mu \left\{ \frac{(1 - L_0/L)_1}{\eta} + \frac{\eta_1 L_0}{2\pi L} \int_0^{2\pi} \psi(\theta) \, d\theta \right\},
\]

(4.3.6)

\[
\tau_w = \frac{\mu \eta_1}{\left[1 - (1 - \mu)\beta^4 \left(R/R_0\right)^4 + 8k \left(R/R_0\right)^2/R_0^3 - 4\sqrt{k} \left(R/R_0\right)^2/\alpha R_0\right]},
\]

(4.3.7)

\[
\tau_s = \left[\tau_w \right]_{R/R_n=1-\delta/R_0},
\]

(4.3.8)
where \( \psi(\theta) = [\varphi(z)]_{R_{n=0}^{r=b}} \), \( a = 1 - \frac{\delta}{2R_0} \), \( b = \frac{\delta}{2R_0} \),
\[
\eta_1 = 1 + 8k/R_0^2 - 4\sqrt{k}/aR_0 \, , \quad \eta = 1 - (1 - \mu)\beta^2 + 8k/R_0^2 - 4\sqrt{k}/aR_0 \, ,
\]
\[
\lambda = \frac{\tilde{\lambda}}{\lambda_0} \, , \quad (\tau_w, \tau_s) = (\tilde{\tau}_w, \tilde{\tau}_s)/\tau_0 \, ,
\]
\[
\lambda_0 = 8\mu L/\eta_1 \pi R_0^4 \quad \text{and} \quad \tau_0 = 4\mu Q/\eta_1 \pi R_0^3 \quad \text{are the flow resistance and shear stress, respectively for a single-layered Newtonian fluid in a normal artery (no stenosis) with permeable wall and} \quad (\tilde{\lambda}, \tilde{\tau}_w, \tilde{\tau}_s) \quad \text{are respectively, (the impedance, the wall shear stress and the shearing stress at stenosis throat) in their non-dimensional form obtained from the definitions} \quad (\text{Young, 1968}; \tilde{\lambda} = \Delta p/Q, \tilde{\tau}_w = -(R/2) dp/dz \quad \text{and} \quad \tilde{\tau}_s = [\tilde{\tau}_w]_{R_{n=0}^{r=R_0}} \, .
\]

### 4.4 Numerical Results and Discussion

To discuss the results of the study quantitatively, computer codes are developed to evaluate the analytical result for flow resistance, \( \lambda \), the wall shear stress, \( \tau_w \), and shear stress at the stenosis throat, \( \tau_s \) obtained above in Eqs. (4.3.6)-(4.3.8) for various parameter values and some of the critical results are displaced graphically in Figs. (4.4.1)-(4.4.12). The various parameters are selected (Young, 1968; Beavers and Joseph, 1967; Srivastava et al., 2012) as: \( L_0 \, (cm) = 1; \quad L \, (cm) = 1, 2, 5, 10; \quad \alpha = 0.1, 0.2, 0.3, 0.5; \quad \sqrt{k} = 0, 0.1, 0.2, 0.3, 0.4, 0.5; \quad \beta = 1, 0.95, 0.90; \quad \mu = 1, 0.5, 0.3, 0.1; \quad \text{and} \quad \delta/R_0 = 0, 0.5, 0.10, 0.45, 0.20; \quad \text{etc.} \) It is worth mentioning here that present study corresponds to impermeable artery case, to single-layered model study, and to no stenosis case for parameter values \( \sqrt{k} \) (here and after called Darcy number) = 0; \( \beta = 1 \) or \( \mu = 1 \), and \( \delta/R_0 = 0 \); respectively.
The flow resistance $\lambda$, increases with the stenosis height, $\delta/R_0$, for any given set of parameters. At any given stenosis height, $\delta/R_0$, $\lambda$ decreases with the peripheral layer viscosity, $\mu$ from its maximal magnitude obtained in a single-layered study (i.e., $\mu=1$ or $\beta=1$; Fig. 4.4.1). One observes that at any given stenosis height, $\delta/R_0$, the impedance, $\lambda$, increases with the slip parameter, $\alpha$ (Fig. 4.4.2). The blood flow characteristic, $\lambda$, increases with the Darcy number, $\sqrt{k}$ at any given stenosis height, $\delta/R_0$ (Fig. 4.4.3). The impedance, $\lambda$, decreases with increasing tube length $L$ which interns implies that $\lambda$ increases with increasing value of $L_0/L$ (stenosis length, Fig. 4.4.4).

**Figure: 4.4.1. Impedance, $\lambda$ versus stenosis height, $\delta/R_0$ for different $\mu$.**
Figure: 4.4.2. Impedance, $\lambda$ versus stenosis height, $\delta/R_o$ for different $\alpha$.

Figure: 4.4.3. Impedance, $\lambda$ versus stenosis height, $\delta/R_o$ for different $k^{1/2}$. 

$L = L_o = 1$
$\beta = 0.95$
$k^{1/2} = 0.1$
$\mu = 0.3$
$\alpha = 0.1$
Numbers $\alpha$
Numbers $k^{1/2}$
One observes that the flow resistance, $\lambda$, decreases rapidly with increasing value of the Darcy number, $\sqrt{k}$ from its maximal magnitude at $\sqrt{k} = 0$ (impermeable wall) in the range $0 \leq \sqrt{k} \leq 0.15$ and afterwards assumes an asymptotic value with increasing values of the Darcy number, $\sqrt{k}$ (Fig. 4.4.5). We notice that the blood flow characteristic, $\lambda$, increases with the slip parameter, $\alpha$ from its minimal magnitude at $\alpha = 0.1$ and approaches to an asymptotic magnitude when $\alpha$ increases from 0.2 (Fig. 4.4.6).
Chapter 4

4.4 Numerical Results and Discussion

Figure: 4.4.5. Impedance, $\lambda$ versus Darcy number, $k^{1/2}$ for different $\delta/R_0$.

Figure: 4.4.6. Impedance, $\lambda$ versus slip parameter, $\alpha$ for different $\delta/R_0$. 

$L=L_0=1$

$\beta=0.95$

$\mu=0.3$

$\alpha=0.1$

Numbers $\delta/R_0$
The wall shear in the stenotic region, $\tau_w$, increases from its approached value at $z/L_0 = 0$ to its peak value at $z/L_0 = 0.5$ and then decreases from its peak value to its approached value at the end point of the constriction profile at $z/L_0 = 1$ for any given set of parameters Figs. (4.4.7)-(4.4.10). The blood flow characteristic, $\tau_w$, decreases with the peripheral layer viscosity, $\mu$ at any axial location of the constriction profile (Fig. 4.4.7). At any point of stenotic region, the wall shear stress, $\tau_w$, increases with Darcy number, $\sqrt{k}$ (Fig. 4.4.8). The flow characteristic, $\tau_w$, also increases with the slip parameter, $\alpha$ at any axial location in the stenotic region (Fig. 4.4.9). For any given set of parameters, the wall shear stress, $\tau_w$, increases with the stenosis height, $\delta/R_0$ (Fig. 4.4.10).

Figure: 4.4.7. Wall shear stress distribution, $\tau_w$, in the stenotic region for different $\mu$. 
Figure: 4.4.8. Wall shear stress distribution, $\tau_w$ in the stenotic region for different $k^{1/2}$.

Figure: 4.4.9. Wall shear stress distribution, $\tau_w$ in the stenotic region for different $\alpha$. 

Chapter 4

4.4 Numerical Results and Discussion
Figure: 4.4.10. Wall shear stress distribution, $\tau_w$ in the stenotic region for different $\delta/R_o$.

Figure: 4.4.11. Shear stress at stenosis throat, $\tau_s$ versus stenosis height, $\delta/R_o$ for different $\mu$. 

Chapter 4

4.4 Numerical Results and Discussion
The blood flow characteristic, $\tau_s$, increases with the slip parameter, $\alpha$ (Fig. 4.4.12) for any given set of other parameters. Numerical results reveal that the variations of the shear stress, $\tau_s$, are similar to that of the impedance (flow resistances), $\lambda$, with respect to any parameter.

4.5 Conclusions

To observe the effects of the permeability of the artery wall and the peripheral layer on blood flow characteristics due to the presence of a stenosis, a two-fluid blood flow of Newtonian fluid through an axisymmetric stenosis in an artery with permeable wall has been studied. The study enables one to observe the simultaneous effects of the wall permeability and the peripheral layer on blood flow characteristics due to the presence of a stenosis. For any given set of parameters, the blood flow characteristics (impedance, wall shear stress, etc.) assume lower magnitude in two-fluid model than
its corresponding value in one-fluid analysis. The impedance decreases with increasing Darcy number from its maximal magnitude in the case of impermeable wall (i.e., at zero Darcy number). It is therefore concluded that the existence of permeability in the artery wall and the presence of the peripheral layer in the artery help the functioning of the diseased artery.