CHAPTER - VIII

8. Growth and Computational Simulation Analysis on the NLO Crystal 3-Pyridinecarboxamide

8.1 Introduction to 3-Pyridinecarboxamide

NLO is concerned with the interaction of electromagnetic fields with various media to produce new electromagnetic fields altered in phase, frequency or amplitude from the incident fields. The investigations on materials with potential for large Second-Order Non Linear Optical (SONLO) applications are very much emphasized in view of the key-role played by them in the frequency conversion process required for the production of new laser output covering ultraviolet and blue-green region. Recently extensive research is going on materials with the capable of efficient frequency conversion of infrared (IR) radiation to visible and ultraviolet (UV) wavelengths. Especially, those materials, which can generate highly efficient second harmonic blue light using laser diodes, are of great interest for applications including optical computing, optical information processing, optical disk data storage, medical diagnostics, etc. Materials with large nonlinear coefficients, and appropriate refractive index (or birefringence) for phase matching, very low absorption, good optical quality, hardness and ability to take high polish, low
two photon absorption coefficient and high laser damage threshold are needed in order to realize many of these applications [1,2]. In the recent past, extensive investigations are being carried out on organic NLO [3-6] materials due to their high SHG efficiency, synthetic flexibility and better laser damage threshold compared with inorganic materials. The origin of nonlinearity in these materials is due to the presence of delocalized electron systems connecting donor and acceptor groups, which enhance their asymmetric polarizability. This report is focused on the growth of crystals of 3-PyridineCarboxamide (3PCA) and also to study the mechanical response of the crystal using Vickers micro hardness technique. The elastic stiffness constant (C_{11}) were calculated for different loads using Wooster’s empirical formula. The gap between highest occupied molecular orbital and the lowest unoccupied molecular orbital is calculated, which is useful for estimating the efficiency of second harmonic generation.

8.2 Solubility Studies

The commercially available 3PCA (C_{6}H_{6}N_{2}O) was further purified by repeated recrystallization process. In order to obtain single crystals of high quality, purification of starting material was found to be an important step. The recrystallized salt was the charge material for the growth of 3PCA. To grow bulk crystals from solution by slow evaporation technique, it is
desirable to select a solvent in which it is moderately soluble. The size of a crystal depends on the amount of material available in the solution, which in turn is decided by the solubility of the material in that solvent. Hence, we have determined the solubility as de ionized water. Solubility in de ionized water was found good and the crystals grown were found to have better shape and transparency. Good transparent single crystals were obtained after a week. Figure 8.1 shows the as grown crystal of 3PCA with an optimized solution pH value of 3.5.

![Figure 8.1 Grown crystal of 3PCA](image)

**8.3 Powder XRD Analysis**

From the XRD data the crystal possess monoclinic symmetry with the space group P21/c, a well-known noncentrosymmetric space group thus satisfying the demands for second-order NLO activity. The lattice parameter values of
the crystal have been calculated using least-squares fit method and they are found to be \( a = 3.931\text{Å}, \ b = 15.607\text{Å}, \ c = 9.472\text{Å}, \ \alpha = \gamma = 90^\circ, \ \beta = 98^\circ \) respectively. The crystallographic data obtained in the present study were found to be apt with the data reported in literature. The chemical structure of 3PCA was depicted in Figure 8.2.

![Chemical Structure of 3PCA](image)

**Figure 8.2 Chemical Structure of 3PCA**

### 8.4 Optical Studies

The optical transmittance spectrum of 2mm thickness polished 3PCA crystal was recorded using Perkin Elmer Lamda 35 spectrophotometer and is shown in Figure 8.3. Normally aromatic compound absorbs light in UV region and the highly conjugated aromatic molecules absorb light in the visible region. The lower cut-off wavelength is found at 220 nm. This is assigned to \( n \) to \( \pi^* \) transition and it may be attributed to the excitation of charge from the aromatic ring also this transition gives white colour to the grown crystal. High transmittances are observed at the Nd: a YAG laser fundamental and second harmonic wavelength which indicates that the grown crystal is
well suited for frequency conversion of 1064 nm laser radiation. The optical transparency characteristic of the grown crystal is an attractive for opto-electronic applications [7] since these compounds do not absorb radiation in the visible region of the electronic spectrum. An absorption peak observed at 190nm is due to the presence of amide group \( \pi - \pi^* \) transition.

![Graph showing transmittance vs. wavelength](image)

**Figure 8.3 UV–Vis spectrum of 3PCA**

### 8.5 Energy Band gap

The variation of \((\alpha h \nu)^2\) versus incident photon energy \((h \nu)\) using the Tauc[8] relation in the fundamental absorption region was plotted and is shown in Figure 8.4. The optical band gap energy \((E_g)\) was estimated by extrapolation of the linear portion of the curve to a point \((\alpha h \nu)^2 = 0\). Using this
method, the optical energy band gap of the crystal is found as 5.8 eV. This high value of band gap energy and the lower cut-off wavelength of the crystal assert the suitability of the crystal for photonic and optical applications [9].

![Graph of (αhν)^2 vs. E=hv](image)

**Figure 8.4 The variation of (αhν)^2 and incident photon energy (hν) of 3PCA**

### 8.6 Micro Hardness Studies

Micro hardness measurements [10-12] were carried out using Leitz Weitzler hardness tester fitted with a diamond indentor. All the indentation measurements were carried out on the freshly cleaved samples. The indentation was made by varying the load from 25 to 200 gm and the time of indentation was kept at 5 sec. The indented impressions were approximately square. The crystal surfaces were indented at different sites. Diagonal
lengths of the indented impression were measured using calibrated micrometer attached to the eyepiece of the microscope. Several indentations were made on each sample. The average value of the diagonal lengths of indentation mark was used to calculate the hardness. The Microhardness is calculated using the expression $H_v = 1.8544P/d^2$ Kg mm$^{-2}$, where $P$ is the applied load in Kg and $d$ the average diagonal length of the Vickers impression in mm after unloading. It is observed that the increase in hardness with applied load. Although hardness has been defined in several ways, it is now generally accepted that it is the resistance offered to dislocation motion.

![Figure 8.5 Plot of $H_v$ versus Load of 3PCA](image-url)
There are several contributions to the resistance to the dislocation motion and they can be classified into two types

(i) The intrinsic resistance which depends on some structure insensitive physical parameter of the crystal 

(ii) A disorder parameter which depends on the concentration of the imperfections.

The non-liner variation of micro hardness is due to the presence of imperfections. These imperfections can be vacancies, impurity-vacancy pairs, dislocations, low angle grain boundaries etc. The plot of hardness \((H_v)\) versus load for the title crystal is shown in Figure 8.5.

**8.7 Elastic Stiffness Constant**

The elastic stiffness constant \((C_{11})\) was calculated for different loads using Wooster’s empirical formula [13] \(C_{11}= [H_v]^{7/4}\). The \(C_{11}\) values are shown in Table 8.1. These values give an idea about the tightness of bonding between neighbouring atoms. The high value shows that the binding forces between the atoms were quite strong. It was calculated for the loads from 25 to 200 gm. The Plot of \(H_v\) versus Elastic Stiffness Constant for the title crystal is shown in Figure 8.6.
Table 8.1 Elastic Stiffness Constant

<table>
<thead>
<tr>
<th>Load (g)</th>
<th>$C_{11} \times 10^{14}\text{Pa}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>25</td>
<td>9.8116</td>
</tr>
<tr>
<td>50</td>
<td>15.043</td>
</tr>
<tr>
<td>100</td>
<td>21.9706</td>
</tr>
<tr>
<td>200</td>
<td>24.4128</td>
</tr>
</tbody>
</table>

8.8 Kurtz Powder SHG Test

A Q-switched Nd: YAG laser beam of wavelength 1064 nm with an input power of 2.5 mJ and pulse width 8 ns with a repetition rate of 10 Hz is used to illuminating the powdered sample. The grown crystal is powdered with a uniform particle size of 125–150 μm, and then packed in a micro capillary of uniform bore and exposed to the laser radiation [14,15]. The
SHG conversion efficiency is found to be 3 times greater value of the KDP (9 mV). This increased value may be due to the presence of amide in the grown material.

### 8.9 Hyperpolarizability Calculation

For calculating the hyperpolarizability, the geometry of the investigated compound is treated as an isolated molecule. The optimization has been carried out in the unrestricted open-shell Hartree-Fock level. The geometries are fully optimized without any constraint with the help of analytical gradient procedure implemented within Gaussian98W program [16]. The electric dipole moment and dispersion free first-order hyperpolarizability are calculated using finite field method. The finite field method offers a straightforward approach to the calculation of hyperpolarizability. The 3-21(d,p) basis set gives remarkably good geometries for such a small basis set and in fact it is used for the geometry optimization of some high accuracy energy methods.

The nonlinear properties of an isolated molecule in an electric field $E_i(\omega)$ can be represented by the Taylor expansion of the total dipole moment $\mu_i$ induced by the field (taken at zero field),

$$\text{Dipole moment } \mu = \begin{bmatrix} E \\ F_i \\ F_j \end{bmatrix} \quad [8.1]$$
Components of polarizability tensor:

\[
\alpha_{ij} = \begin{bmatrix}
E \\
F_i
\end{bmatrix}
\]  \[8.2\]

Components of hyperpolarizability tensor:

\[
\beta_{ijk} \begin{bmatrix}
\cdot \\
\cdot \\
\cdot
\end{bmatrix}
\]  \[8.3\]

These components are to be distorted by an external electric field. The value of total static polarizability and hyperpolarizability are obtained from the following equation,

\[
\beta_{ijk} = \beta_{ij} + \beta_{ji} + \beta_{jk} + \beta_{kj} + \beta_{ik} + \beta_{ki} + \beta_{il} + \beta_{li} \]  \[8.4\]

In the presence of an applied electric field, first order hyperpolarizability is a third rank tensor that can be described by a 3×3×3 matrix. The components of the 3D matrix can be reduced to 10 components because of the Kleinman symmetry [17]. The matrix can be given in the lower tetrahedral format. It is obvious that the lower part of the 3×3×3 matrix is a tetrahedral. The calculation of NLO properties with high accuracy is challenging and requires consideration of many different issues. Computational techniques are becoming valuable in designing, modeling and screening novel NLO materials. The calculated value of hyperpolarizability for the title compound is \(0.818586 \times 10^{-30}\) esu, which is nearly 4 times that of urea \((0.20920 \times 10^{-30}\) esu). This highest value of hyperpolarizability may be due to the presence of electro-
negative nitro group and π bonds. The calculated values of dipole moment and hyperpolarizability values are tabulated in Table 8.2. The $\beta_{zzz}$ direction shows biggest value of hyperpolarizability which insists that the delocalization of electron cloud is more that direction than other directions. The NLO responses can be understood by examining the energetic of frontier molecular orbitals.

Table 8.2 The dipole moment ($\mu$) and first-order hyperpolarizibility ($\beta$) of Glutamic acid derived from DFT calculation

<table>
<thead>
<tr>
<th>$\beta_{xxx}$</th>
<th>0.204492</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\beta_{xyy}$</td>
<td>0.0565447</td>
</tr>
<tr>
<td>$\beta_{xyy}$</td>
<td>-0.013891</td>
</tr>
<tr>
<td>$\beta_{yyyy}$</td>
<td>-0.0026034</td>
</tr>
<tr>
<td>$\beta_{zxx}$</td>
<td>-0.75063</td>
</tr>
<tr>
<td>$\beta_{xyz}$</td>
<td>0.105752</td>
</tr>
<tr>
<td>$\beta_{zyy}$</td>
<td>-0.17586</td>
</tr>
<tr>
<td>$\beta_{xzz}$</td>
<td>-0.21523</td>
</tr>
<tr>
<td>$\beta_{yzz}$</td>
<td>0.0221795</td>
</tr>
<tr>
<td>$\beta_{zzz}$</td>
<td>0.111822</td>
</tr>
<tr>
<td>$\beta_{tot}$</td>
<td>0.818586</td>
</tr>
<tr>
<td>$\mu_x$</td>
<td>1.491223686</td>
</tr>
<tr>
<td>$\mu_y$</td>
<td>0.0000627771</td>
</tr>
<tr>
<td>$\mu_z$</td>
<td>2.227837147</td>
</tr>
<tr>
<td>$\mu$</td>
<td>1.928502945</td>
</tr>
</tbody>
</table>
Figure 8.7(a) and Figure 8.7(b) shows the HOMO and LUMO of 3PCA. There is an inverse relationship between hyperpolarizability and HOMO–LUMO.

HOMO energy = $-0.298$ a.u

LUMO energy = $0.208$ a.u

HOMO–LUMO energy gap = $0.506$ a.u

Figure 8.7[a, b] HOMO and LUMO representation of 3PCA
8.10 Electronic Excitation Mechanism

The static polarizability value [18,19] is proportional to the optical intensity and inversely proportional to the cube of transition energy. With this concept, larger oscillator strength ($f_n$) and $\Delta \mu_{gn}$ with lower transition energy ($E_{gn}$) is favourable to obtain large first static polarizability values. Electronic excitation energies, oscillator strength and nature of the respective excited states were calculated by the closed-shell singlet calculation method and are summarized in Table 8.3.

**Table 8.3 Computed absorption wavelength ($\lambda_{ng}$), energy ($E_{ng}$), oscillator strength ($f_n$) and its major contribution.**

<table>
<thead>
<tr>
<th>n</th>
<th>$\lambda_{ng}$ (nm)</th>
<th>$E_{ng}$ (eV)</th>
<th>$f_n$</th>
<th>Major contribution</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>228.2</td>
<td>5.43</td>
<td>0.0009</td>
<td>H-4-&gt;L+1(30%), H-3-&gt;L+1(21%), H-4-&gt;L+14(9%), H-4-&gt;L+19(5%)</td>
</tr>
<tr>
<td>2</td>
<td>199.4</td>
<td>6.22</td>
<td>0.1948</td>
<td>H-0-&gt;L+1(71%), H-1-&gt;L+4(10%)</td>
</tr>
<tr>
<td>3</td>
<td>197.6</td>
<td>6.27</td>
<td>0.0147</td>
<td>H-2-&gt;L+4(32%), H-2-&gt;L+1(11%), H-3-&gt;L+4(11%), H-2-&gt;L+16(6%)</td>
</tr>
</tbody>
</table>

*(Assignment; H=HOMO, L=LUMO, L+1=LUMO+1, etc.)*

The orbital involved in the electronic transition for different values of HOMO, LUMO are represented in Figure 8.8.
Figure 8.8 Representation of the orbital involved in the electronic transition for (a) HOMO−4 (b) LUMO+1 (c) LUMO+14 (d) HOMO-0 (e) HOMO-2 (f) LUMO+4 (g) HOMO-3 (h) HOMO-1 (i) LUMO+16
8.11 Conclusion

Good quality crystals of 3PCA were grown from aqueous solution by slow evaporation technique at room temperature. The optical transmittance window, lower cut off wavelength, band gap of the title crystal has been identified by UV-VIS-NIR studies. The results of the response of the mechanical behaviour of the crystal will have significant effect on machining the crystal for device purpose. DFT calculations were performed for the title NLO crystal to evaluate the first order hyperpolarizability value. The calculated value of hyperpolarizability for the title compound is $0.818586 \times 10^{-30}$ esu, which is nearly 4 times that of urea ($0.20920 \times 10^{-30}$ esu). This could be a new series of prospective crystals for second-order NLO applications such as electro optic or frequency doubler of conventional diode lasers.
References


